# Low-temperature sintering mechanism on uranium dioxide

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Abstract Based on a point defect model, the mechanism of low-temperature sintering of uranium dioxide was studied in this paper. The diffusion coefficient of uranium in  $UO_{2+x}$ , sintering temperature and densification equation in low-temperature sintering were analyzed by both the point defect model and low-temperature sintering experiments. The results showed that the diffusion activation energy of uranium in over-stoichiometric  $UO_{2+x}$  was lowered by 3.0 ev than that in stoichiometric  $UO<sub>2</sub>$ . And the diffusion coefficient of uranium in  $UO_{2+x}$  was proportional to  $x^2$ . In addition, the theoretical low temperature sintering temperature was calculated to be in the range of 1089–1151 °C, which indicated that it was necessary to maintain proper over-stoichiometric oxygen for low-temperature sintering process. Moreover, the calculation results by the point defect model matched perfectly with the experiment results, PDM might be a good model to describe the relationships between defects concentration and atmosphere composition.

## Introduction

Recently, high-temperature sintering kinetics of  $UO_{2+x}$  [[1,](#page-4-0) [2](#page-4-0)], diffusion coefficient of oxygen and uranium and diffusion equation [[3,](#page-4-0) [4](#page-4-0)] have been widely studied. Belle [[5\]](#page-4-0) considered that the model of sintering  $UO<sub>2</sub>$  consisted of

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three parts: shrinkage of volume, spheroidization of void, growth of crystal. Since then many researchers have investigated low-temperature sintering of  $UO<sub>2</sub>$  pellets [\[6](#page-4-0)– [8](#page-4-0)], however the low-temperature sintering mechanism was rarely reported. In this study, we investigated the lowtemperature mechanism, feasibility and influencing factors of  $UO<sub>2</sub>$  pellets. We established a relationship between diffusion coefficient, pyknosis speed and partial pressure of oxygen (Po<sub>2</sub> or x). Based on this relationship, we could optimize some technological parameters for lower-temperature sintering of  $UO<sub>2</sub>$  pellets.

## Volume diffusion coefficients of uranium ions in PDM

When sintering, uranium ions diffused by vacancy in  $UO<sub>2</sub>$ . The diffusion behavior was described by the following equation:

$$
D_U = D_V[V_U'''''] \tag{2.1}
$$

where  $D_V$  was the diffusion coefficient of uranium vacancy, and  $D_U$  was the diffusion coefficient of uranium.

To simplify the discussion, we made the following assumptions:  $\odot$  uranium vacancies diffused accurately by a three-dimensional random walk process in  $UO<sub>2</sub>$  crystals, and the transition direction was random,  $\oslash$  only volume diffusion was considered; surface diffusion and interface diffusion were neglected,  $\circled{3}$  the uranium vacancies transmitted in  $UO<sub>2</sub>$  by exchanging positions with the nearest uranium ion only,  $\Phi$  the lattice parameters of  $UO<sub>2</sub>$  crystals were considered to be constant when interstitial oxygen anions were dissolved in the crystals.

We assumed a component gradient existed along the z-axis in one-dimensional random walk process, and

vacancies moved a distance of  $\lambda$  at each jump. The adjacent lattice surfaces were apart in a distance of  $\lambda$ , and the numbers of solute vacancies passing through two adjacent surfaces in unit area were  $n_1$  and  $n_2$ , respectively. We also assumed the transition frequency,  $v$ , was the average value of each vacancy jumping out from the plane surfaces per second. Therefore, the vacancy flux of moving from plane 1 to plane 2 in unit time was:

$$
J = \frac{n_1 - n_2}{2}v
$$
 (2.2)

And the component gradient in unit distance was:

$$
\frac{n_1 - n_2}{\lambda^2} = -\frac{\partial c}{\partial z} \tag{2.3}
$$

From (2.2) and (2.3), we obtained

$$
\mathbf{J} = -\frac{1}{2}\lambda^2 v \frac{\partial c}{\partial z} \tag{2.4}
$$

From the first diffusion law,  $J = -D \frac{\partial c}{\partial z}$  (2.5), we got the diffusion coefficient of one-dimensional random walk process

$$
D = \frac{1}{2}\lambda^2 v \tag{2.5}
$$

So the diffusion coefficient of three-dimensional random walk process was  $1/3$  of the D in  $(2.5)$  as we have assumed above. Note that the sub lattice structure of uranium ion is face-centered cubic, and we also know that the geometry parameter, f, was 0.781 [\[9](#page-4-0)]. Based on the above analyses, the following equation was derived to describe the diffusion coefficients of three-dimensional random walk process:

$$
D_V = \frac{1}{6} f \lambda^2 v = 0.1302 \lambda^2 v \tag{2.6}
$$

We assumed that the jump frequency of uranium vacancy was  $v_0$  and the energy barrier of transition was  $\Delta G_m$ , then the diffusion coefficient of uranium ion was

$$
D_U = 1.5624\lambda^2 v_0 \exp\left(-\frac{\Delta G_m}{kT}\right)[V_U'''''] \tag{2.7}
$$

In combination with the results in literature  $[10]$  $[10]$ , we got the equations underneath.

When the value of  $Po_2$  or x was very small,

$$
D_{Ux} \approx D_{U0} = 1.5624 \lambda^2 v_0 \exp\left(-\frac{\Delta G_m + \Delta G_S - \Delta G_{Fo}}{kT}\right)
$$
\n(2.8)

However, when the value of  $Po<sub>2</sub>$  or x was quite large or very large

$$
D_{Ux} = 1.5624 \lambda^2 v_0 \exp\left(-\frac{\Delta G_m + \Delta G_S - 2\Delta G_{Fo}}{kT}\right) x^2 \tag{2.9}
$$

where  $D_{U0}$  was the diffusion coefficient of uranium ion in stoichiometry  $UO_2$ ,  $D_{Ux}$  was the diffusion coefficient of uranium ion in over-stoichiometry  $UO_{2+x}$ . From (2.8) to (2.9), we know that the volume diffusion coefficient of uranium ion was related to temperature only when  $Po<sub>2</sub>$  or x was very small indicating the diffusion process of uranium ion was an intrinsic diffusion. At a certain temperature, the volume diffusion coefficient of uranium ion increased with the increment of Po<sub>2</sub> or x, then the diffusion process of uranium ion became an extrinsic diffusion. These findings were consistent with previous experimental results.

Note that  $D_0$  and Q (activation energy of diffusion) can be expressed with the following equations.

When  $Po_2$  or x was very small,

$$
D_{00} = 1.5624 \lambda^2 v_0 \tag{2.10}
$$

$$
Q_0 = \Delta G_m + \Delta G_S - \Delta G_{Fo}
$$
 (2.11)

When  $Po_2$  or x was quite large,

$$
D_{0x} = 1.5624\lambda^2 v_0 x^2 \tag{2.12}
$$

$$
Q_x = \Delta G_m + \Delta G_S - 2\Delta G_{Fo}
$$
\n(2.13)

In the equations,  $Q_0$  was the activation energy of diffusion in stoichiometry  $UO_2$ , and  $Q_X$  was the activation energy of diffusion in over-stoichiometry  $UO_{2+x}$ .

From (2.10) to (2.13) we found it was feasible to sinter uranium dioxide in low-temperature if the  $UO_{2+x}$  pellets were used. And there was a relationship between the diffusion activation energies of uranium ion:

$$
Q_x - Q_0 = -\Delta G_{\text{Fo}} \tag{2.15}
$$

From  $(2.14)$  we found that the diffusion activation energy of one uranium ion in  $UO_{2+x}$ ,  $Q_x$  was lower than that in  $UO_2$ ,  $Q_0$ . The difference in Gibbs free energy  $G_{F_0}$  was 3.0 ev. These findings were in good agreement with those reported by Matzke [[11\]](#page-4-0) if the diffusion activation entropy was ignored. When temperature was constant, the diffusion activation energy, Q, can be expressed as

$$
Q = \Delta H - T\Delta S \approx \Delta H \tag{2.15}
$$

The diffusion activation entropy of uranium ion in  $UO_{2+x}$ ,  $\Delta$  H<sub>Ux</sub> was about 1.8–3.0 ev according to the

<span id="page-2-0"></span>**Table 1** Values of at 1500  $^{\circ}$ C

X	0.0045	0.033	0.08	0.1	
$D_{Ux}$ (other reports)	$4.45 \times 10^{-14}$ (Lindner)	$8.15 \times 10^{-14}$ (Hawkins)	$5.33 \times 10^{-13}$ (Matzke)	$7.04 \times 10^{-13}$ (Matzke)	
$D_{Ux}$ (this paper)	$3.77 \times 10^{-18}$	$2.03 \times 10^{-16}$	$1.19 \times 10^{-15}$	$1.86 \times 10^{-15}$	
$D_{\text{Ux}}^{\text{b}}$ (this paper)	$1.15 \times 10^{-15}$	$6.19 \times 10^{-14}$	$3.64 \times 10^{-13}$	$5.68 \times 10^{-13}$	
$D_{\text{Ux}}^s$ (this paper)	$3.03 \times 10^{-14}$	$1.63 \times 10^{-12}$	$9.56 \times 10^{-12}$	$1.49 \times 10^{-11}$	

experimental results in literature [[3–7\]](#page-4-0). Therefore, the diffusion activation energies of uranium ion were obtained as follows:  $Q_x \approx 2.5$  ev and  $Q_0 \approx 5.5$  ev. So, when Po<sub>2</sub> or x was quite large, the diffusion coefficient of uranium ion was

$$
D_{Ux} = 1.5624 \lambda^2 v_0 x^2 \exp\left(-\frac{2.5}{kT}\right)
$$
 (2.16)

Because the lattice constant of  $UO<sub>2</sub>$  was  $5.474 \times 10^{-7}$ <sup>10</sup> m [\[12](#page-4-0)], and unidirectional transition frequency  $(v_0)$  was about  $10^{-13}/s$  in solid. We obtained

$$
D_{Ux} = 2.341 \times 10^{-6} x^2 \exp\left(-\frac{2.5}{kT}\right)
$$
 (2.17)

When the temperature was  $1500 \degree C$ ,

$$
lg D_{Ux} = -12.73 + 2lg x \tag{2.18}
$$

The above equation was almost the same with the experimental finding by Matzke [[11\]](#page-4-0) and Marin and Con-tamin [\[13](#page-4-0)]. Table 1 shows a comparison of  $D_{Ux}$  between the calculation and literatures. The results indicated that we could reduce the sintering temperature of uranium dioxide pellets by adding  $U_3O_8$  powder that changed  $U_2$  to  $UO_{2+x}$ .

# The temperature of low-temperature sintering uranium dioxide

Let  $T'$  and  $T''$  represent the sintering temperature of hightemperature sintering process and low-temperature sintering process, respectively. When  $D_{U0} = D_{Ux}$ , we had the following relationships between  $T''$  and x or Po<sub>2</sub>

$$
\frac{Q_0 - \Delta G_{\text{Fo}}}{T''} = 2k \ln x + \frac{Q_0}{T'}
$$
\n(3.1)

$$
\frac{Q_0 - \Delta G_{\text{Fo}}}{T''} = 0.03633k + \frac{1}{3}k \ln P_{\text{O}_2} + \frac{Q_0}{T'}
$$
(3.2)

From  $(3.1)$  to  $(3.2)$ , the theoretical sintering temperature of low-temperature sintering process was calculated to be in the range of 1089–1151 °C when  $x = 0.25$ . This was in good agreement with the experimental value.

## The sintering kinetics of low-temperature sintering uranium dioxide

The uranium vacancy transferred from neck surface to adjacent spherical face in a volume diffusion process. The following equations were used to describe this process.

$$
\frac{dV}{dt} = J Sa^3 \tag{4.1}
$$

$$
S = 2\pi b \rho = \frac{\pi b^3}{r}
$$
 (4.2)

$$
V = \pi b^2 \rho = \frac{\pi b^4}{2r} \tag{4.3}
$$

$$
J = D_{Ux} \frac{\Delta C_V}{\rho} \tag{4.4}
$$

where  $dt$  was the time parameter,  $dV$  was the volume parameter, J was the number of uranium vacancies that leave the neck in an unit area and unit time, S was the area of neck, and a was the radius of uranium atoms.

Considered the interfacial force of powder was  $\sigma$  that was constant to certain pellets. The concentration of uranium vacancies was  $[V''_{U}]$  based on the PDM calculations. According to the literature [\[14](#page-4-0)], we know that the difference of concentrations between uranium vacancies that form in neck surface and grain surface was given by the equation:



<span id="page-3-0"></span>
$$
\Delta C_V = \frac{\sigma a^3 [V_U'''']}{kT}
$$
\n(4.5)

Substituting (4.5) to Sintering Neck Model then followed by transpositions, we obtained the equation of neck diffusion kinetics:

$$
\left(\frac{\mathbf{b}}{\mathbf{r}}\right)^5 = \frac{40\sigma a^3 \mathbf{D}_{\mathbf{U}x}}{\mathbf{r}^3 \mathbf{k} \mathbf{T}} \mathbf{t} \tag{4.6}
$$

This equation matched well with the result reported by Kingery, Huang [\[15](#page-4-0)] and Johnson. From (4.6), we know the growth speed of neck increased when the value of x increased. So over-stoichiometry of oxygens in uranium dioxide was necessary to achieve low temperature sintering of UO<sub>2</sub> pellets.

Assumed there were powder grains with K in mass that accumulated to form a pore after sintering. Thus the sintering density was:



Fig. 1 The appearances of sintering pellets in weak-oxidizingatmosphere-sintering

$$
\rho_{t} = \frac{K\rho_{th}}{K + 8\left(1 - \frac{b}{r}\right)^{3}}\tag{4.7}
$$

where  $\rho_{\text{th}}$  and t were the theoretical density and sintering time, respectively.

Table 3 Results of oxidizing-atmosphere-sintering

Time	Temperature				
	800 $\mathrm{^{\circ}C}$	$1000 \text{ °C}$	1200 °C	1400 °C	
1 <sub>h</sub>	$O/U = 2.155$ $\rho_1 = 6.851(62.50\%)$	$O/U = 2.118$ $\rho_1 = 9.171(83.68\%)$ $d_1 = 8.679$ $e = 11.93\%$	$O/U = 2.054$ $\rho_1 = 9.727(88.75\%)$ $d_1 = 8.497$ $e = 5.63\%$	$O/U = 2.033$ $\rho_1 = 10.221(93.26\%)$ $d_1 = 8.380$ $e = 0.30\%$	
2 <sub>h</sub>	$O/U = 2.135$ $\rho_2$ = 7.253 (66.18%)	$O/U = 2.022$ $\rho_2$ = 9.292 (87.78%) $d_2 = 8.660$ $e = 10.33\%$	$O/U = 2.022$ $\rho_2 = 9.807(89.48\%)$ $d_2 = 8.460$ $e = 3.45\%$	$O/U = 2.048$ $\rho_2 = 10.251(93.53\%)$ $d_2 = 8.369$ $e = 0.50\%$	
3 <sub>h</sub>	$O/U = 2.096$ $\rho_3 = 7.716$ (70.40%) $d_3 = 9.237$ $e = 26.25\%$	$O/U = 2.023$ $\rho_3 = 9.377(85.56\%)$ $d_3 = 8.628$ $e = 10.02\%$	$O/U = 2.014$ $\rho_3 = 9.900 (90.33\%)$ $d_3 = 8.427$ $e = 2.424\%$	$O/U = 2.047$ $\rho_3 = 10.299(93.97\%)$ $d_3 = 8.359$ $e = 0.40\%$	

Fig. 2 The micro-structures of sintering pellets in weakoxidizing-atmosphere-sintering



Sintering techniques Theoretical value **Experimental** value Error value  $(%$  $1200 °C$ C  $2-3 h$   $0.850$   $0.899$   $5.5$ 1400 °C C  $2-3 h$   $0.850$   $0.933$   $9.3$ 

<span id="page-4-0"></span>Table 4 Comparisons of theoretical linear shrinkage and testing datas

#### Experimental and model checking out

Materials and methods

The main properties of original powder are listed in the Table [2](#page-2-0). The additive pellets were produced by industrial technologies. The additives include  $U_3O_8$  (15.11 wt.%), ammonium oxalate (0.20 wt.%), and zinc stearate  $(0.10 \text{ wt.}\%)$ . The original diameter,  $d_0$  and original average density,  $\rho_0$  of pellets were 14.00 mm and 5.825 g/cm [3], respectively.

The sintering experiments were carried out in a vertical molybdenum wire furnace which provided the highest temperature of  $1600 \pm 1$  °C. High-purity N<sub>2</sub> (99.99%), high-purity  $H_2$  (99.99%) and pure  $CO_2$  (99.9%) were used. The oxidizing atmosphere was composed of  $N_2 + 1.5\%$ CO2. The samples were heated at a heating rate of less than 400  $\degree$ C/h, then were held in the temperature range of 600–  $650$  °C for 0.5 h. Subsequently they were cooled at a cooling rate of less than 400  $\degree$ C/h.

The density and porosity of sintered pellets were examined by GB11927-89. O/U was measured by GB11842-89. And the grain size and diameter of pellets were measured by GB/T6394-2002 and micrometer, respectively.

#### Results and verification

The results of low-temperature sintering process experiments are summarized in Table [3.](#page-3-0) The appearance and microstructure of the as-produced pellets are shown in Fig. [1](#page-3-0) and [2](#page-3-0), respectively. From Table [3,](#page-3-0) we know that the equation [16],  $x = 1.018 \exp(0.017/T) \text{Po}_2^{\frac{1}{6}}$ , was correct.

The calculation results of linear contraction are listed in Table 4. We can see that the difference between theoretical and experimental values was less than 10%.

#### **Conclusions**

Low-temperature sintering mechanism of uranium dioxide was for the first time studied by PDM. The essence of lowtemperature sintering of uranium dioxide was activated sintering. The results showed that the diffusion activation energy for uranium in over-stoichiometric  $UO_{2+x}$  was lowered by 3.0 ev than that in stoichiometric  $UO_2$ , and  $D_{Ux}$  $\propto$  x<sup>2</sup>. Thus it was necessary to maintain proper over-stoichiometric oxygen for low-temperature sintering process.

The diffusion coefficient of uranium ion in  $UO_{2+x}$  can be expressed by the equation  $D_{Ux} = 1.5624 \lambda^2 v_0 x^2$  $exp(-(2.5/kT))$ , based on volume diffusion mechanism, the equations of growth sintering neck and densification kinetics were  $(b^5/r^2) = (60\sigma a^3D_{Ux}/kT)t$  and  $\rho_t = Kd_{th}/$  $(K + 8(1 - (b/r<sup>3</sup>)))$ , respectively.

The theoretical sintering temperature calculated by PDM was in the range of  $1089-1151$  °C, which perfectly matched with the experimental results.

The PDM ignored the effect of association of defects, but it could describe the relationships of defects concentration and atmosphere quite well. The findings in this study indicate that PDM could be an effective tool for investigating low-temperature sintering mechanism of uranium dioxide. We also predict that the PDM might be an effective tool for investigating low-temperature sintering mechanisms of other non-stoichiometric compounds. It might be helpful to decrease the sintering temperature when people manufacture industrial products.

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